DETERMINATION OF AIR-TIGHTNESS OF THE PACKAGINGS OF ELECTRONIC DEVICES BY THE THERMOACOUSTIC METHOD

M. MALIŃSKI
Technical University of Koszalin
Department of Electronics
Śniadeckich 2, 75-328 Koszalin, Poland
email: mmalin@tu.koszalin.pl

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The method of determination of air-tightness of packagings based on the measurement of the thermoacoustic signals generated by transistor chips is described in this paper. This method makes possible the estimation of the radius of the hole in the packaging and, as a result, its air-tightness. In this paper the fitting of the theoretical model to experimental results is presented and discussed. Because of its simplicity and nondestructive character, the presented method can be applied in the quality control departments of the electronic industry.

Key words: air-tightness, thermoacoustics.

1. Introduction

Air-tightness of the packaging is an important parameter which must be controlled in the electronic industry works. The measuring techniques are very often troublesome because they require either high pressures and special liquids, as it is in the case of a pressure bomb, or gases as it is in the trace gas method [1, 2], or elevated temperatures and liquids [3] or low pressure and liquids [3], depending on the kind of a test. Most often these methods are time consuming and also destructive. The aim of this work was to test and estimate the threshold level of the air leaks that can be determined by the thermoacoustic technique. For the purpose of interpretation of the thermoacoustic results, the electric model of a gas transport from the packaging to the thermoacoustic cell is presented and discussed in this paper. The method of detection presented in the paper is especially useful for metal packagings TO-3, TO-18, TO-39, TO-72 of electronic devices and for the packagings of hybrid circuits. The idea of the testing method presented in this paper is based on the theory of a thermophone elaborated and published by H.D. ARNOLD and I. B. CRANDALL [4]. Thermophones were used in the past for pressure calibration of measuring microphones since it was expected that it was possible to calculate this way the absolute value of acoustic pressure in the closed chamber.
and its frequency characteristics. The modern mathematical model of both the photo-
and thermoacoustic effect was elaborated by A. ROSENCAIG and A. GERSHO in 1976 [5]. The theory of the thermal-piston and mechanical-piston effects responsible for the
thermoacoustic effect, used for the detection of air-tightness described in this paper, was
next elaborated by H. HU et al. in 1999 [6]. The results of investigations of air-tightness
of the plastic packaging of power transistors with a photoacoustic microscopic thermal
wave method with the optical heat flux generation were presented by Z. SUSZYŃSKI
in papers [7, 8]. In this method, the areas of exfoliation of the plastic encapsulant from
the metal lead frames were localized. Similar method of thermal wave imaging with a
thermoacoustic detection was also applied for the detection of a place of exfoliation of
the packaging by A. BODERMAN in 1996 [9] and A. ROSENCAIG and J. OPSAL in
1986 [10].

2. Theory

The flow of the mass of gas through the pipe of radius \( r \) and length \( L \) can be
described by the Poiseuille Eq. (1):

\[
\frac{dm}{t} = \frac{\rho \cdot \pi \cdot r^4}{\eta \cdot 8 \cdot L} \cdot dp,
\]

where \( dm \) – mass of the gas, \( t \) – time of flow, \( \rho \) – density of the gas, \( \eta \) – viscosity of gas,
\( L \) – length of the pipe or thickness of metal wall of the packaging, \( r \) – is the radius of
the hole. By analogue with the Ohm law \( I = V/R \), the resistance of the gas flow can be
determined as:

\[
R = \frac{\eta \cdot 8 \cdot L}{\pi \cdot r^4 \cdot \rho} = \frac{dp \cdot t}{dm}.
\]

Because of the Clapeyron equation, this amount of gas \( dm \) flowing into the chamber of
volume \( V \) causes the increase of the overpressure \( dp \) given by the Eq. (3).

\[
dp = \frac{dm \cdot N_a \cdot k \cdot T}{M \cdot V}.
\]

Here \( N_a \) is the Avogadro number, \( k \) – Boltzman constant, \( T \) – temperature, \( V \) – volume
of the chamber, \( M \) – molar mass of the air.

Thus the heat capacity of the chamber can be determined as:

\[
C_i = \frac{M \cdot V_i}{N_a \cdot k \cdot T}.
\]

It is an analogue of the electric condenser where \( V = Q/C \) and \( Q \) is the electric
charge, \( V \) is the voltage and \( C \) is the capacity of the condenser.

The schematic diagram of the thermoacoustic chamber with the transistor in the
metal case is presented in Fig. 1.
Fig. 1. A schematic diagram of the thermoacoustic chamber with a microphone. Description: 1 – microphone, 2 – thermoacoustic chamber, 3 – transistor, 4 – stopper.

The idea of the presented thermoacoustic (TA) method is based on the periodic heat generation in volume $V_1$ of a transistor (3) and measurement of the periodic overpressure in the volume $V_2$ of the thermoacoustic chamber (2), being result of the air flow from volume $V_1$ to $V_2$ through the hole in the packaging of a transistor. Frequency amplitude and phase TA characteristics bring information about the radius of the hole and in general, about the air tightness of the packaging.

Changes of the pressure in the thermoacoustic chamber can be described in the model being the electric analog of the gas flow from the volume of the transistor to the volume of the thermoacoustic chamber. In this model pressure corresponds to voltage, electric capacity to heat capacity, resistance of the mass flow to electric resistance, velocity of mass transport to electric current. The schematic diagram of the electric model describing the gas flow is presented in Fig. 2.

![Fig. 2. An electric model of the gas flow. $R$ is the resistance, $C_i$ ($i = 1, 2$) are capacities. $G$ is the current generator.](image)

The value of the voltage $U_2(R)$ is given by Eq. (5).

$$ U_2(R) = \left[ \frac{-i}{\omega(C_1 + C_2) + i\omega^2C_1C_2R} \right] \cdot I. $$  (5)
The value of $U_2(R = 0)$ is given by Formula (6).  

$$U_2(R = 0) = \left[ \frac{-i}{\omega(C_1 + C_2)} \right] \cdot I.$$  

(6)

Thus the ratio of voltages is given by Formula (7). It expresses the ratio of the voltages and the corresponding pressures of the packaging exhibiting a hole, expressed by the resistance $R$ and the open packaging ($R = 0$) corresponding the case of a huge hole in the packaging.

$$\frac{U_2(R)}{U_2(0)} = \frac{\omega(C_1 + C_2)}{\omega(C_1 + C_2) + i\omega^2C_1C_2R}.$$  

(7)

The modulus of the ratio of the voltages (overpressures) and the phase delay of the thermoacoustic signal relative to the electric current $I$ are given by Formula (8) and (9) respectively when $C_i$ and $R$ are given by Eqs. (2) and (4).

$$\left| \frac{U_2(r)}{U_2(0)} \right| = \frac{1}{\sqrt{1 + (\omega\tau(r))^2}},$$  

(8)

$$\varphi = \arctg(-\omega\tau(r)) \frac{180}{\pi},$$  

(9)

$$\tau(r) = \frac{\eta 8LMV_2}{\pi \cdot r^4 \rho N_a k T (1 + V_2/V_1)}.$$  

(10)

For the computations presented below the following values of the parameters were taken:

- $\eta = 1.7 \cdot 10^{-5} [\text{Ns/m}^2]$, $L = 10^{-4} [\text{m}]$, $M = 28 \cdot 10^{-3} [\text{kg/mole}]$,
- $V_2 = 2.5 \cdot 10^{-6} [\text{m}^3]$, $V_2/V_1 = 5$, $r = 20 \mu\text{m} ... 60 \mu\text{m}$,
- $\rho = 1.3 [\text{kg/m}^3]$, $N_a = 6 \cdot 10^{23} [\text{mole}^{-1}]$,
- $T = 300 \text{ K}$, $k = 1.38 \cdot 10^{-23} [\text{J/K}]$.

The results of theoretical computations of the amplitude of the thermoacoustic (TA) signal versus the frequency of modulation are presented in Fig. 3.

The results of theoretical computations of the phase of the TA signal versus the frequency of modulation are presented in Fig. 4.

Because the calibration characteristics of the system are often not known and the TA signal depends, among others, on the thickness of a transistor chip and the quality of soldering of the chip to the lead frame [11–13], the measurements of the relative change of the amplitude of the TA signal with the frequency of modulation and the relative phase shift of the TA signals are recommended. The theoretical curves of the $|U_2(r)/U_2(0)|$ versus the frequency of modulation are presented in Fig. 5.

Theoretical curves of the relative phase shift of the TA signals $U_2(r)$ and $U_2(0)$ versus the frequency of modulation are presented in Fig. 6.
Fig. 3. Amplitudes of the TA signal versus the frequency of modulation for different values of the radius of the hole in the transistor packaging: Lines: 1) $r = 500 \ \mu m$ (open transistor), 2) $r = 60 \ \mu m$, 3) $r = 50 \ \mu m$, 4) $r = 40 \ \mu m$, 5) $r = 30 \ \mu m$.

Fig. 4. Phases of the TA signal versus frequency of modulation for different values of a radius of the hole. Lines: 1) $r = 500 \ \mu m$ (open transistor), 2) $r = 60 \ \mu m$, 3) $r = 50 \ \mu m$, 4) $r = 40 \ \mu m$, 5) $r = 30 \ \mu m$. 
Fig. 5. Ratio of the amplitudes $|U_2(r)/U_2(0)|$ of the thermoacoustic signals versus the frequency of modulation for different values of the radius of the hole: 1) $r = 500 \, \mu m$ (open transistor), 2) $r = 60 \, \mu m$, 3) $r = 50 \, \mu m$, 4) $r = 40 \, \mu m$, 5) $r = 30 \, \mu m$.

Fig. 6. Relative phase shifts of the thermoacoustic signals $U_2(r)$ and $U_2(0)$ versus the frequency of modulation for different values of the radius of the hole: 1) $r = 500 \, \mu m$ (open transistor), 2) $r = 60 \, \mu m$, 3) $r = 50 \, \mu m$, 4) $r = 40 \, \mu m$, 5) $r = 30 \, \mu m$. 
The electric model enables also the analysis of the dependences of the amplitude ratio of the TA signals on the radius of the hole for a series of modulation frequencies. The theoretical curves are presented in Fig. 7.

The dependences of the phase shift of TA signals on the radius of the hole for a series of modulation frequencies are presented in Fig. 8.

Fig. 7. Dependence of the amplitude ratio \( \frac{U_2(r)}{U_2(0)} \) on the radius of the hole for three frequencies of modulation. Description of lines: 1) \( f = 5 \) Hz, 2) \( f = 20 \) Hz, 3) \( f = 80 \) Hz.

Fig. 8. Dependence of the relative phase shift of the thermoacoustic signals \( U_2(r) \) and \( U_2(0) \) on the radius of the hole for three frequencies of modulation. Description of lines: 1) \( f = 5 \) Hz, 2) \( f = 20 \) Hz, 3) \( f = 80 \) Hz.
3. Experimental results

Transistors in the common collector circuits were supplied with the constant voltage $U_{CE} = 15$ V and their collector currents $I_C$ were controlled with a sinusoidal voltage generator in the base-emitter circuit $U_{BE}$. Collector currents were periodically modulated in the range $I_C = 10 \text{ mA} - 0 \text{ mA}$ so that the maximum instantaneous electric power dissipated in the transistor was $P = 150 \text{ mW}$. Thermoacoustic signals in the thermoacoustic chamber were detected by the electret microphone with a preamplifier. For the maximum power $P = 150 \text{ mW}$ and the frequency of modulation $f = 20 \text{ Hz}$, the amplitude of the thermoacoustic signal was $U_2 = 200 \text{ mV}$ for an open transistor while for a closed transistor in the same conditions $U_2 = 0 \text{ mV}$. The phase was measured between the TA signal and the voltage on the emitter resistor $U_{RE}$ proportional to the collector current $I_C$. It was measured both with a digital oscilloscope and the electronic phase-meter. Amplitude of a TA signal was also measured with the digital oscilloscope and after the rectifier – with a digital multimeter.

For the purpose of testing of the proposed model, holes were done mechanically in the metal packagings of several transistors and a correlation between the theoretical model and experimental results was investigated.

Fittings of theoretical amplitude and phase curves to experimental results, obtained for the packaging TO39 of transistors BF178, are shown in Figs. 9 and 10 respectively.

Fitting of theoretical curves of the amplitude ratio and phase shift to experimental characteristics obtained for two other transistors is presented in Figs. 11 and 12 respectively.

![Fig. 9. Fitting of theoretical curves of the amplitude ratio to experimental results obtained for two transistors exhibiting different values of the radius of holes. Description of lines: 1) – open transistor, 2) $r = 60 \mu\text{m}$, 3) $r = 50 \mu\text{m}$, 4) $r = 42 \mu\text{m}$, 5) $r = 32 \mu\text{m}$.]
Fig. 10. Fitting of theoretical curves of the relative phase shift to experimental results obtained for two transistors exhibiting different values of the radius of holes. Description of lines: 1) – open transistor, 2) \( r = 60 \, \mu m \), 3) \( r = 50 \, \mu m \), 4) \( r = 45 \, \mu m \), 5) \( r = 34 \, \mu m \).

Fig. 11. Fitting of theoretical curves of the amplitude ratio to experimental results obtained for two transistors exhibiting different values of the radius of holes. Description of lines: 1) – open transistor, 2) \( r = 60 \, \mu m \), 3) \( r = 50 \, \mu m \), 4) \( r = 34 \, \mu m \), 5) \( r = 29 \, \mu m \).
Fig. 12. Fitting of theoretical curves of the amplitude ratio to experimental results obtained for two transistors exhibiting different values of the radius of holes. Description of lines: 1) – open transistor, 2) $r = 60 \, \text{µm}$, 3) $r = 50 \, \text{µm}$, 4) $r = 34 \, \text{µm}$, 5) $r = 29 \, \text{µm}$.

4. Conclusions

The results of investigations of the correlation of the amplitude and phase thermoacoustic characteristics obtained for a series of transistor samples, exhibiting holes in the packaging, indicate that it is possible to determine the air tightness of packagings of electronic devices by this method. The measurements also present information about the sensitivity of this thermoacoustic approach. A good correlation of experimental results and theoretical characteristics, concerning both amplitude and phase, proves the correctness of the presented electric model of gas flow from the packaging to the thermoacoustic chamber.

References


